

## Equation of state from complex Langevin simulations

Attanasio, Felipe; Jäger, Benjamin; Ziegler, Felix P.G.

*Published in:*  
EPJ Web of Conferences

*DOI:*  
10.1051/epjconf/202227405012

*Publication date:*  
2022

*Document version:*  
Final published version

*Document license:*  
CC BY

*Citation for pulished version (APA):*  
Attanasio, F., Jäger, B., & Ziegler, F. P. G. (2022). Equation of state from complex Langevin simulations. *EPJ Web of Conferences*, 274, [05012]. <https://doi.org/10.1051/epjconf/202227405012>

Go to publication entry in University of Southern Denmark's Research Portal

### Terms of use

This work is brought to you by the University of Southern Denmark.  
Unless otherwise specified it has been shared according to the terms for self-archiving.  
If no other license is stated, these terms apply:

- You may download this work for personal use only.
- You may not further distribute the material or use it for any profit-making activity or commercial gain
- You may freely distribute the URL identifying this open access version

If you believe that this document breaches copyright please contact us providing details and we will investigate your claim.  
Please direct all enquiries to [puresupport@bib.sdu.dk](mailto:puresupport@bib.sdu.dk)

# Equation of state from complex Langevin simulations

Felipe Attanasio<sup>1,\*</sup>, Benjamin Jäger<sup>2,\*\*</sup>, and Felix P.G. Ziegler<sup>3,\*\*\*</sup>

<sup>1</sup>Institute for Theoretical Physics, Universität Heidelberg, Philosophenweg 16, D-69120, Germany

<sup>2</sup>CP3-Origins & Danish IAS, Department of Mathematics and Computer Science, University of Southern Denmark, Campusvej 55, 5230 Odense M, Denmark

<sup>3</sup>School of Physics and Astronomy, The University of Edinburgh, EH9 3FD Edinburgh, United Kingdom

**Abstract.** We use complex Langevin simulations to study the QCD phase diagram with two light quark flavours. In this study, we use Wilson fermions with an intermediate pion mass of  $\sim 480$  MeV. By studying thermodynamic quantities, in particular at lower temperatures, we are able to describe the equation of state.

## 1 Introduction

Determining the QCD phase diagram from first principles is a long-standing problem and has implications for heavy-ion collision experiments and cosmology. Standard lattice methods provide a non-perturbative approach for QCD, which has reached an impressive sub-percent precision. The main obstacle to direct lattice determination of the QCD phase diagram is the numerical *sign* problem. For a non-vanishing chemical potential, the fermion determinant is complex, which prohibits standard importance sampling based Monte Carlo methods to be used. In recent years, a wide range of methods has been proposed to overcome or circumvent the sign problem in QCD. However, the best approach has not yet been identified. Here we present our results using the complex Langevin method [1, 2], which is based on stochastic quantisation. Over recent years, a large amount of improvements have been proposed and developed [3–12], which have made the complex Langevin method a potential contender for determining the QCD phase diagram. We report on our findings attempting an early determination of the equation of state at an intermediate pion mass.

## 2 Setup

We use a lattice setup with two mass-degenerate dynamical quarks. The lattice parameters are: the hopping parameter is  $\kappa = 0.1544$  and the gauge coupling is  $\beta = 5.8$ . This implies a fixed-scale approach. These parameters are based on the simulations of [13]. Our determination of the lattice spacing, performed via the Wilson flow, has led to  $a \sim 0.06$  fm, consistent with the [13]. We use a spatial volume of  $24^3$  in lattice units. To map the phase diagram, we vary the temporal extent from 4 up to 128 sites, but most of our simulations are in the region

\*e-mail: [pyfelipe@thphys.uni-heidelberg.de](mailto:pyfelipe@thphys.uni-heidelberg.de)

\*\*e-mail: [jaeger@imada.sdu.dk](mailto:jaeger@imada.sdu.dk)

\*\*\*e-mail: [felix.ziegler@ed.ac.uk](mailto:felix.ziegler@ed.ac.uk)

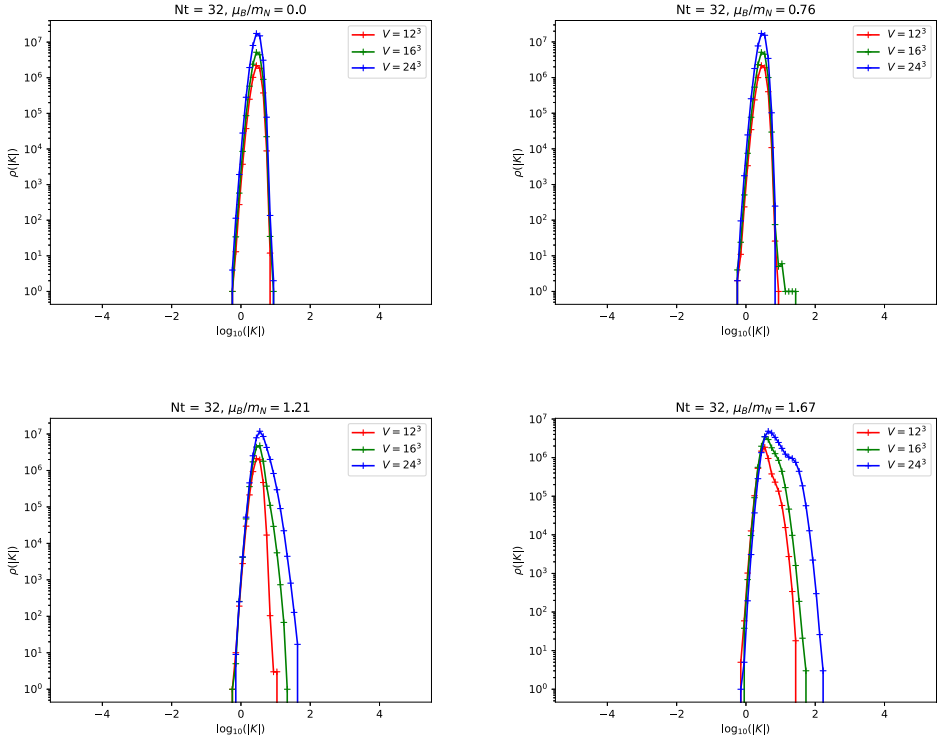


Figure 1: Histograms of the Langevin force  $K$  for  $N_t = 32$  ( $T = 103$  MeV) for three volumes and different values of the baryon chemical potential  $\mu_B$  in units of the nucleon mass.

of  $N_t \leq 32$ . The temperature range covers 822 to 25 MeV. For the quark chemical potential, we use values  $a\mu = 0-2$ , in physical units up to  $\sim 6000$  MeV. To overcome runaway solutions and instabilities, we use adaptive step size scaling. In addition, we used one step of gauge cooling after every Langevin step. For this setup, we combined it with dynamic stabilisation (DS) with  $\alpha_{DS} = 1$ . Here we employ a modified version of DS for which we do not use the adaptive step size scaling in the additional DS drift; instead, we keep the step size fixed in order to ensure that the additional force term only depends on the distance to SU(3), i.e.,

$$\bar{\epsilon} \alpha_{DS} K_{DS}. \quad (1)$$

The step size of the dynamic stabilization force is fixed to  $\bar{\epsilon} = 10^{-3}$ . For our simulations, we use a modified version of the openQCD v1.6 [14] and openQCD-fastsum software packages [15]. As an indication of the stability of the simulation, we show histograms of the Langevin force for various baryon chemical potentials<sup>1</sup>, which can be seen in Figure 1. These results were also obtained at smaller volumes, i.e.  $16^3$  and  $12^3$ . We find that all histograms show a localised distribution, which typically is taken as an indication for correct convergence of the CL method [7]. With increasing  $\mu$  the histograms show some structure emerging, which indicates that the chemical potential is changing the Langevin forces. However, the histograms or skirts remain sufficiently compact over the entire range.

<sup>1</sup>The baryon chemical potential  $\mu_B$  is related to the quark chemical potential  $\mu$  by:  $\mu = \mu_B/3$ .

### 3 Results

To identify the phase boundaries, we study the Polyakov loop and the fermion density. Fig-

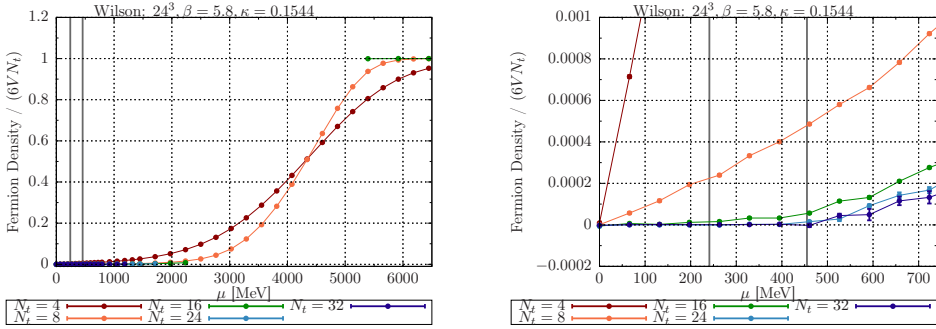


Figure 2: Fermion density as a function of the quark chemical potential. The vertical grey lines correspond to  $m_\pi/2$  and  $m_N/3$  respectively.

Figure 2 shows the fermion density as a function of the quark chemical potential. The plot on the left shows the entire simulation range. We found that the density saturates as expected for a large enough density. This is a lattice artifact that appears if all lattice sites are filled with  $N_{\text{spin}} \times N_c = 6$  fermionic degrees of freedom per quark flavour. Adding more is not possible because of Pauli blocking. On the right hand side of Figure 2 we show an enlarged view, which highlights the region up to 700 MeV. As expected, colder simulations show a later rise of the density. For the lowest temperature shown here, we see that density remains approximately zero until  $m_N/3$ , where we expect that quarks can be created. This phenomenon is also known as the Silver Blaze and emerges dynamically. Figure 3 shows the Polyakov loop. On the left-hand side as a function of the chemical potential. The same lattice artefact as seen in the density is visible. However, here it causes the Polyakov loop to shrink again. The right panel of Figure 3 provides a different perspective on the transition. For larger chemical potentials, the transition appears to occur at lower temperatures. This also is in line with our expectations that the transition temperature reduces as the chemical potential increases. We

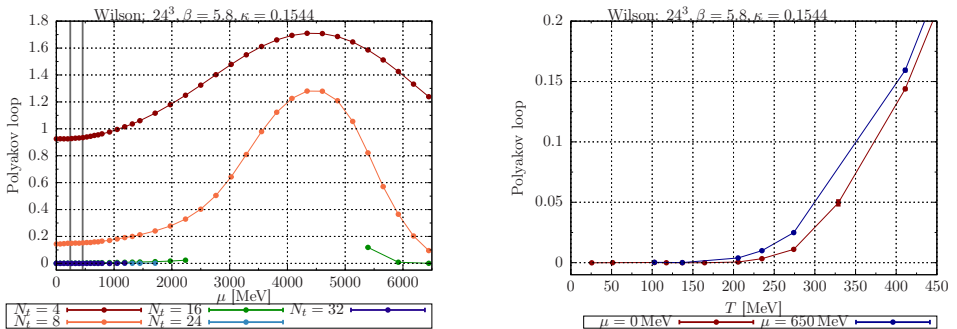


Figure 3: Polyakov loop as a function of the quark chemical potential (left) and temperature (right). The saturation density is reached at  $\mu \sim 4000$  MeV.

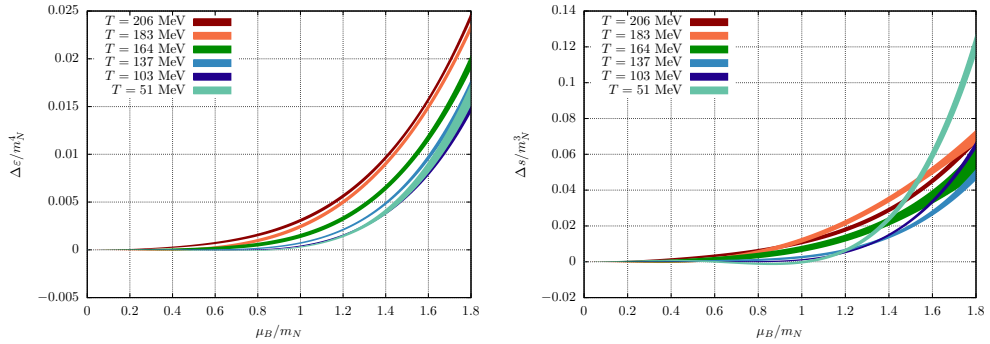


Figure 4: The equation of state extracted from the fermion density.

fit the fermion density to a cubic polynomial and use the result to compute the pressure and from it the equation of state. Figure 4 shows the energy and entropy density, respectively. We find that the equation of state becomes stiffer with decreasing temperatures.

## 4 Conclusion and Outlook

Here we presented an update on our efforts to determine the equation of state from complex Langevin simulations. Although our simulations are still at an unphysical quark mass, which corresponds to a pion mass of roughly 480 MeV, our simulations confirm the expectations on the QCD phase diagram. In the future, we will improve our simulation by using  $O(a)$  improved actions. We hope that this allows us also to reduce the pion mass closer to the physical world. In addition, we plan to study multiple volumes to get a handle on the order of the transition(s).

## 5 Acknowledgement

Computing resources were provided by the University of Southern Denmark, Deic initiative (UCloud, GenomeDK) and PRACE on Hawk (Stuttgart) with project ID 2018194714. This work used the DiRAC Extreme Scaling service at the University of Edinburgh, operated by the Edinburgh Parallel Computing Centre on behalf of the STFC DiRAC HPC Facility ([www.dirac.ac.uk](http://www.dirac.ac.uk)). We are thankful for the support of the Deutsche Forschungsgemeinschaft (DFG) under EXC2181/1 - 390900948 (STRUCTURES) and the SFB 1225 (ISOQUANT).

## References

- [1] G. Parisi, Y.s. Wu, *Sci. Sin.* **24**, 483 (1981)
- [2] G. Parisi, *Phys. Lett.* **B131**, 393 (1983)
- [3] G. Aarts, I.O. Stamatescu, *JHEP* **09**, 018 (2008), [0807.1597](https://arxiv.org/abs/0807.1597)
- [4] E. Seiler, D. Sexty, I.O. Stamatescu, *Phys. Lett.* **B723**, 213 (2013), [1211.3709](https://arxiv.org/abs/1211.3709)
- [5] K. Nagata, J. Nishimura, S. Shimasaki, *PTEP* **2016**, 013B01 (2016), [1508.02377](https://arxiv.org/abs/1508.02377)
- [6] J. Nishimura, S. Shimasaki, *Phys. Rev.* **D92**, 011501 (2015), [1504.08359](https://arxiv.org/abs/1504.08359)
- [7] K. Nagata, J. Nishimura, S. Shimasaki, *Phys. Rev.* **D94**, 114515 (2016), [1606.07627](https://arxiv.org/abs/1606.07627)
- [8] G. Aarts, E. Seiler, D. Sexty, I.O. Stamatescu, *JHEP* **05**, 044 (2017), [1701.02322](https://arxiv.org/abs/1701.02322)

- [9] F. Attanasio, B. Jäger, Eur. Phys. J. **C79**, 16 (2019), 1808.04400
- [10] K. Nagata, J. Nishimura, S. Shimasaki, JHEP **05**, 004 (2018), 1802.01876
- [11] M. Scherzer, E. Seiler, D. Sexty, I.O. Stamatescu, Phys. Rev. **D101**, 014501 (2020), 1910.09427
- [12] E. Seiler (2020), 2006.04714
- [13] L. Del Debbio, L. Giusti, M. Luscher, R. Petronzio, N. Tantalo, JHEP **02**, 011 (2006), hep-lat/0512021
- [14] <https://luscher.web.cern.ch/luscher/openQCD/> (2020)
- [15] <https://fastsum.gitlab.io/> (2020)